Viscous Flowfield Calculations on Pointed Bodies at Angle of Attack in Nonuniform Freestreams

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Theme

THREE-DIMENSIONAL laminar boundary layers are calculated on cones, at incidence in supersonic wake-like freestreams. The nonuniformity of the freestream makes the flowfield nonconical, and therefore both the inviscid and viscous flowfields are three dimensional.

Content

A wake-like nonuniform freestream is represented by the relation

$$V_1/V_{\infty} = 1 + A[1 - \exp(-By_1^2)] \tag{1}$$

where V_1 is the velocity at the transverse position y_1 , V_{∞} is the velocity at $y_1 = 0$ (nose of the cone), A is the velocity defect at $y_1 = 0$, and B is a constant which controls the distribution of the nonuniformity (see Fig. 1). Other properties in the freestream are

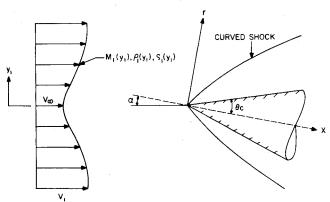


Fig. 1 Coordinate system and freestream properties.

determined from Eq. (1) and the condition of constant static pressure and total enthalpy. The inviscid flowfield is then calculated from a modified form of Rakich's three-dimensional method of characteristics (NASA TN D-5341). Typical results for the non-dimensional pressure $(\bar{p}_b=p_b/\rho_\infty V_\infty^2)$ and inviscid surface Mach Number (M_b) are shown in Fig. 2 for a 20° half-angle cone at $\alpha=10^\circ$ and $M_\infty=3$. For a velocity defect of A=0.2 the surface Mach number goes to unity at x=0.95 (based on B=10) in the windward plane, and the three-dimensional method of characteristics cannot be continued beyond this point. Member AIAA.

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Index category: Boundary Layer and Convective Heat Transfer-

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Three-dimensional laminar boundary layers are calculated by using the axisymmetric analog or small cross-flow assumption. In order to apply the axisymmetric analog, inviscid surface streamlines and the scale factor, or metric coefficient, h_3 must be calculated first. The scale factor h_3 corresponds to the surface coordinate σ which is orthogonal to the streamlines, and h_3 represents the "equivalent radius." With Φ defined as the circumferential angle on the cone and ϕ the angle between the streamline direction and the generator on the cone, the equation of a streamline becomes

$$D\Phi/Dx = \tan \phi/x \sin \theta_c \tag{2}$$

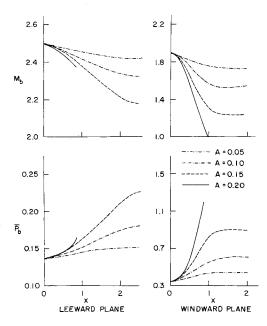


Fig. 2 Surface Mach number and pressure distributions on a cone in axisymmetric wakes; $M_{\infty}=3,\,\theta_{c}=20^{\circ},\,\alpha=10.$

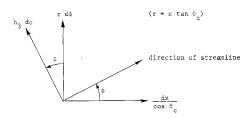


Fig. 3 Streamline coordinates.

where D/Dx is a derivative along a streamline ($\sigma = \text{constant}$), the differential distance orthogonal to the streamlines is $h_3 d\sigma$, which is also the spacing between adjacent streamlines.

From Fig. 3 it follows that

$$h_3 d\sigma = r d\Phi \cos \phi - (dx/\cos \theta_c)\sin \phi \tag{3}$$

and with $\Phi = \Phi(\sigma, x)$ Eq. (3) gives

$$(\partial \Phi/\partial \sigma)x = (h_3/r\cos\phi) \tag{4}$$

By interchanging the order of differentiation, the following equation is formed.

$$\frac{D}{Dx} \left(\frac{\partial \Phi}{\partial \sigma} \right)_{x} = \left[\frac{\partial}{\partial \sigma} \left(\frac{D\Phi}{Dx} \right) \right]_{x} = \left[\frac{\partial}{\partial \Phi} \left(\frac{D\Phi}{Dx} \right) \right]_{x} \left(\frac{\partial \Phi}{\partial \sigma} \right)_{x} \tag{5}$$

Finally, substitute Eqs. (2) and (4) into Eq. (5) to obtain the differential equation for the scale factor as

$$\frac{D}{Dx} \left[\ln \left(\frac{h_3}{r \cos \phi} \right) \right] = \frac{1}{x \sin \theta} \left[\frac{\partial}{\partial \Phi} (\tan \phi) \right]_{x}$$
 (6)

Eqs. (2) and (6) are integrated numerically to determine the angular position Φ and the scale factor h_3 , respectively, along a streamline. The streamline-direction angle ϕ is obtained from the three-dimensional method of characteristics solution on the surface.

Using the axisymmetric analog, laminar boundary-layer properties are calculated along each inviscid surface streamline, independently of the others, by Blottner's numerical method. For the special case of a cone at incidence in a uniform freestream, the present method was found to calculate skin-friction coefficients (C_f) and heating rates (q_w) reasonably close to both experimental data and an "exact" numerical solution.

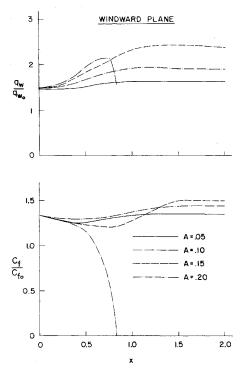


Fig. 4 Normalized local heat transfer and skin-friction coefficient on a cone; $M_{\infty}=3, \, \theta_c=20^{\circ}, \, \alpha=10^{\circ}, \, R_{\rm n}/{\rm ft}=2.85\times 10^5, \, T_{\rm w}/T_o=0.4.$

Figure 4 shows the distribution of q_w/q_{w_o} and C_f/C_{f_o} along the windward plane for the same conditions as the cone in Fig. 2 (C_{f_o} and q_{w_o} refer to values for this cone at $\alpha = 0$ in a uniform

freestream). In Fig. 2 the surface Mach number was found to go to unity at x=0.95 for A=0.2, and Fig. 4 indicates boundary-layer separation at x=0.8 (upstream of the $M_b=1$ position). However, the separation point is affected by the ratio of wall to stagnation temperature (T_w/T_o) , and separation did not occur upstream of the $M_b=1$ position for cold walls, i.e. $T_w/T_o<0.2$. On the other hand, separation was always found to occur upstream of the $M_b=1$ position for cool to hot walls, i.e. $T_w/T_o>0.4$. The Reynolds number per foot (R_n/ft) did not affect the separation point.

Figure 5 depicts the circumferential variation of q_w/q_{w_o} and C_f/C_{f_o} at x = 0.75 (based on B = 10) for several angles of attack.

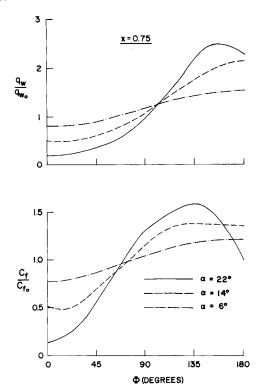


Fig. 5 Normalized local heat transfer and skin-friction coefficient on a cone; $M_{\infty}=3, \theta_{c}=20^{\circ}, \alpha=14^{\circ}, A=0.1, R_{n}/\mathrm{ft}=2.85\times10^{5}.$

The windward plane is $\Phi = 180^{\circ}$, and the leeward plane is $\Phi = 0$. For $\alpha = 22^{\circ}$ both q_w and C_f reach a maximum around the cone away from the windward plane.

In conclusion, the axisymmetric analog for three-dimensional laminar boundary-layers was found to yield accurate q_w and C_f on pointed cones at incidence in uniform freestreams. Calculations for cones at incidence in wake-like nonuniform freestreams indicated an adverse pressure gradient on the surface which reduced the inviscid Mach number to unity in some cases. In addition, boundary-layer separation was found to occur in some examples.

Reference

¹ Blottner, F. G., "Finite Difference Methods of Solutions of the Boundary-Layer Equations," *AIAA Journal*, Vol. 8, No. 2, Feb. 1970, pp. 193–205.